



Zustandsdichten

effektive Masse, Beweglichkeit

Seniorensseminar 6. 12. 2024

Density of states for electrons $D_e(\varepsilon_e)$

definition of states from uncertainty principle $(\Delta x)^3 (\Delta p)^3 \geq h^3$

volume in momentum space for 1 state $(\Delta p)^3 = \frac{h^3}{\Delta x^3} = \frac{h^3}{V_{state}}$

number of momentum states with $|p'| \leq |p|$ $N(|p|) = \frac{4\pi/3 |p|^3}{h^3/V}$
i.e. in a sphere with radius p

We need $N(\varepsilon)$, because occupation depends on energy, not on momentum

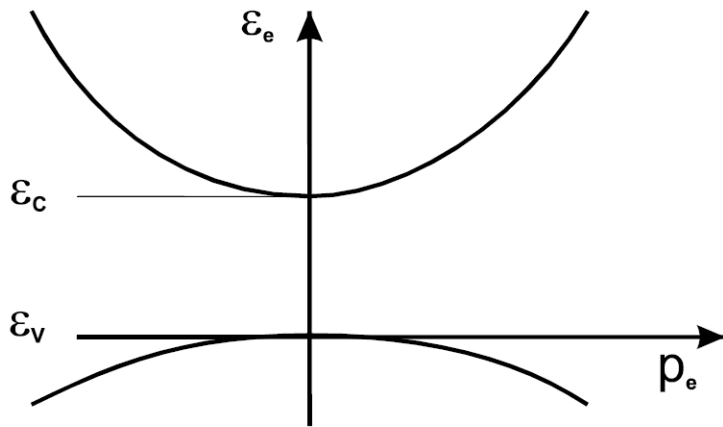
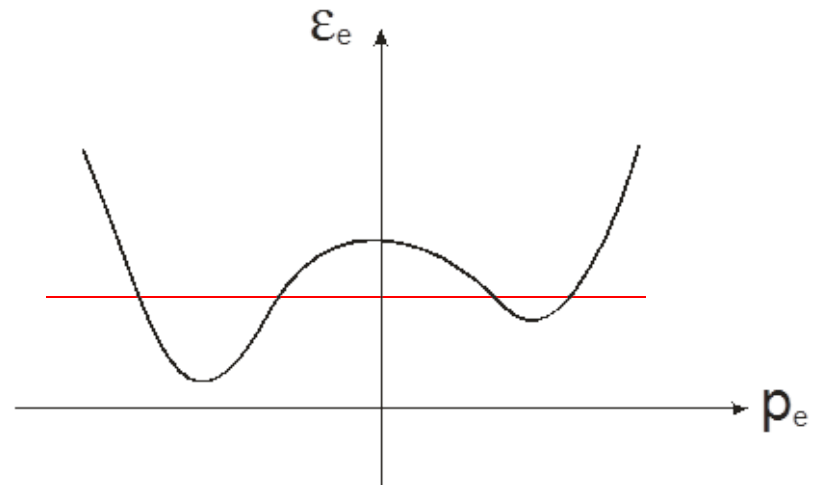
$\varepsilon_e(p) = \varepsilon_C + \alpha p^2 + \dots$ in analogy with free particles $\varepsilon_e - \varepsilon_C = \alpha p^2 = \frac{p^2}{2m_e^*}$

$$N_e(\varepsilon_e) = 2 \overset{\text{spin}}{\downarrow} * \frac{4\pi V (2m_e^*)^{3/2}}{3h^3} (\varepsilon_e - \varepsilon_C)^{3/2}$$

$$D_e(\varepsilon_e) = \frac{1}{V} \frac{dN_e(\varepsilon_e)}{d\varepsilon_e} = 4\pi \left(\frac{2m_e^*}{h^2} \right)^{3/2} (\varepsilon_e - \varepsilon_C)^{1/2}$$

effective mass

can $\epsilon_e(\vec{p}_e)$ be asymmetric?



free electrons

$$\epsilon_e - \epsilon_c = \alpha p_e^2 + \cancel{\beta p_e^4} \dots$$

$$\epsilon = \frac{p^2}{2m} \quad \alpha = \frac{1}{2m_e^*}$$

$$\frac{1}{m_e^*} \approx \frac{d^2 \epsilon_e}{dp_e^2}$$

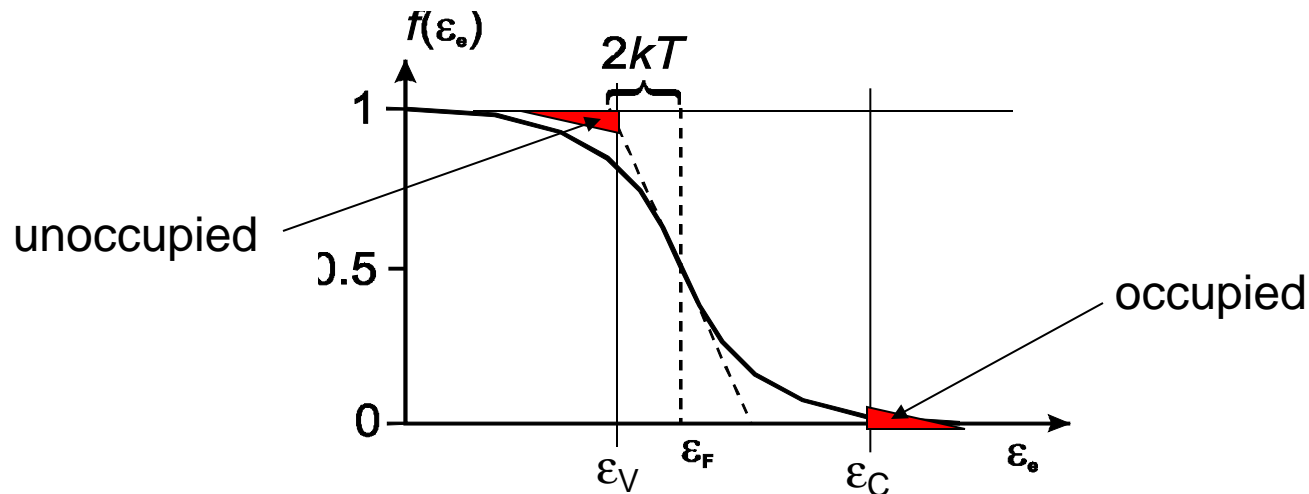
$$\frac{1}{m_e^*} < 0 \quad \text{at} \quad \epsilon_e \approx \epsilon_v$$

Energy distribution of electrons n_e

$$dn_e(\varepsilon_e) = D_e(\varepsilon_e) f_e(\varepsilon_e) d\varepsilon_e$$

density of states $D_e(\varepsilon_e)$ occupation function $f(\varepsilon_e) = \frac{1}{\exp\left(\frac{\varepsilon_e - \varepsilon_F}{kT}\right) + 1}$

principle: minimum of Free Energy $F = E - TS$



condition: temperature T and volume V must be fixed

Electron and hole concentrations in the dark

$$f(\varepsilon) \approx \exp\left(-\frac{\varepsilon - \varepsilon_F}{kT}\right) \quad \text{for } \varepsilon - \varepsilon_F \gg kT$$

electrons

$$n_e = \int_{\varepsilon_C}^{\infty} D(\varepsilon) f(\varepsilon) d\varepsilon = N_C \exp\left(-\frac{\varepsilon_C - \varepsilon_F}{kT}\right); \quad N_C = 2 \left(\frac{2\pi m_e^* kT}{h^2}\right)^{3/2} \quad \text{effective density of states}$$

$$\text{for } m_e^* = m_e \text{ and } T = 300\text{K: } N_C = 2 \times 10^{19} / \text{cm}^3$$

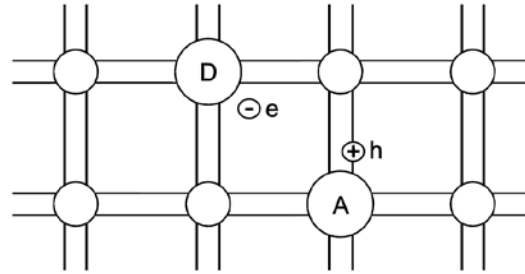
holes

$$n_h = \int_{-\infty}^{\varepsilon_V} D(\varepsilon) [1 - f(\varepsilon)] d\varepsilon = N_V \exp\left(-\frac{\varepsilon_F - \varepsilon_V}{kT}\right) \quad N_V = 2 \left(\frac{2\pi m_h^* kT}{h^2}\right)^{3/2}$$

$$n_e n_h = N_C N_V \exp\left(-\frac{\varepsilon_G}{kT}\right) = n_i^2$$

independent of ε_F
independent of doping

Doping



n – type, donors

p – type, acceptors

1 or more valence electrons more than host

1 or more valence electrons less than host

e.g. P, As in Si
Si on Ga site in GaAs

e.g. In, B in Si
Si on As site in GaAs

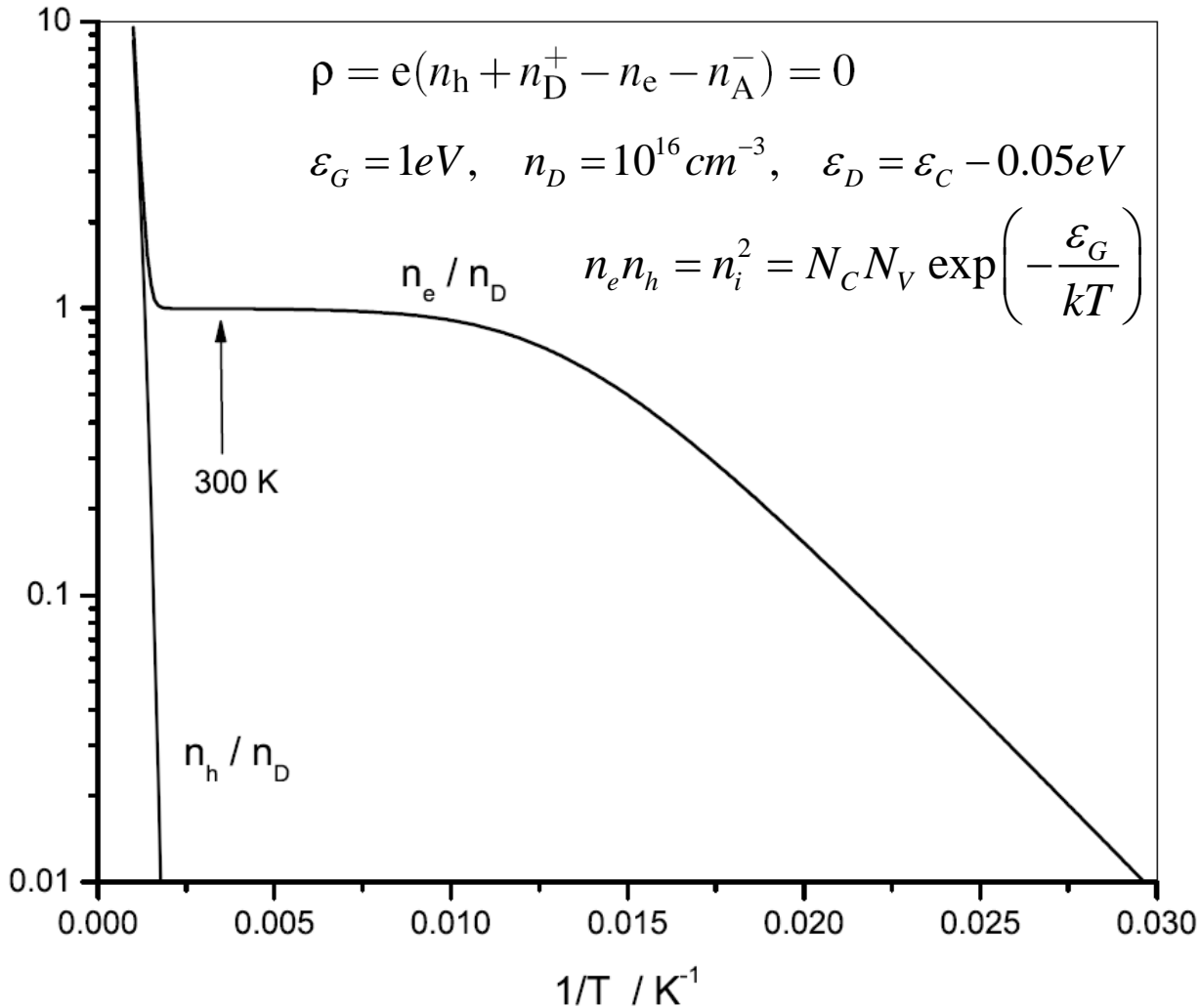
binding energy of electron to donor (hole to acceptor) from hydrogen model

$$\varepsilon_H = \frac{m_e e^4}{2(4\pi\varepsilon_0)^2 \hbar^2} = 13.6 \text{ eV}$$

$$\varepsilon_C - \varepsilon_{e,D} = \varepsilon_H \frac{m_e^*}{m_e \varepsilon^2}$$

order of kT at 300K

Electrons in n-type semiconductor



$$n_h = N_V \exp\left(\frac{\epsilon_V - \epsilon_F}{kT}\right)$$

$$n_D^+ = n_D \left[1 - \frac{1}{\exp\left(\frac{\epsilon_D - \epsilon_F}{kT}\right) + 1} \right]$$

$$n_e = N_C \exp\left(-\frac{\epsilon_C - \epsilon_F}{kT}\right)$$

$$n_A^- = n_A \left[\frac{1}{\exp\left(\frac{\epsilon_A - \epsilon_F}{kT}\right) + 1} \right]$$

Transport properties, drift current

caused by electric field E

acceleration $a_i = \frac{z_i e E}{m_i^*}$

drift current

$$\langle v_i \rangle = \int_0^\infty a_i \exp(-t/\tau_{c,i}) dt = a_i \tau_{c,i}$$

$$j_{Q,f,i} = z_i^2 e n_i b_i E = \sigma_i E$$

$$\langle v_i \rangle = z_i \frac{e}{m_i^*} \tau_{c,i} E .$$

$$j_{Q,f,i} = -\frac{\sigma_i}{z_i e} \text{grad}(z_i e \varphi)$$

mobility $b_i = \frac{e \tau_{c,i}}{m_i^*}$

Diffusion current (looking for a driving force)

diffusion current

$$j_{Q,d,i} = z_i e (-D_i \text{grad } n_i) \quad \text{Fick's law of diffusion}$$

$$j_{Q,d,i} = -z_i e n_i D_i \frac{\text{grad } n_i}{n_i} = -z_i e n_i D_i \text{grad} \left\{ \ln(n_i) \right\}$$

$$\mu_i = \mu_{i,0} + kT \ln \left(\frac{n_i}{N_i} \right)$$

chemical potential of particles i
= free energy of particles i, if uncharged

$$\text{Einstein relation} \quad \frac{b_i}{D_i} = \frac{e}{kT}$$

$$j_{Q,d,i} = -\frac{z_i e n_i b_i}{e} \text{grad } \mu_i = -\frac{\sigma_i}{z_i e} \text{grad } \mu_i$$

total charge current

$$j_{Q,i} = -\frac{\sigma_i}{z_i e} \{ \text{grad} \mu_i + \text{grad}(z_i e \varphi) \}$$

$$j_{Q,i} = -\frac{\sigma_i}{z_i e} \text{grad}(\mu_i + z_i e \varphi) = -\frac{\sigma_i}{z_i e} \text{grad} \eta_i$$

electrochemical potential

$$\eta_i = \mu_i + z_i e \varphi$$

free energy of particles with charge $z_i e$

for electrons ($z_i = -1$) and holes ($z_i = 1$)

$$j_Q = \frac{\sigma_e}{e} \text{grad} \eta_e - \frac{\sigma_h}{e} \text{grad} \eta_h$$

$$\eta_e = \mu_e - e \varphi = \varepsilon_{FC}$$

$$\eta_h = \mu_h + e \varphi = -\varepsilon_{FV}$$

$$j_Q = \frac{\sigma_e}{e} \text{grad} \varepsilon_{FC} + \frac{\sigma_h}{e} \text{grad} \varepsilon_{FV}$$

electron current + hole current

Energies of electron-hole pairs

energy

$$\varepsilon_e + \varepsilon_h = \varepsilon_C + 3/2kT + (-\varepsilon_V) + 3/2kT = \varepsilon_G + 3kT$$

Free Energy: electrochemical energy, electrical energy + chemical energy

$$\eta_e + \eta_h = \mu_e - e\varphi + \mu_h + e\varphi = \mu_e + \mu_h$$

in the dark

$$\varepsilon_{FC} = \varepsilon_{FV} = \varepsilon_F$$

$$\eta_e + \eta_h = \varepsilon_{FC} - \varepsilon_{FV} = \mu_e - e\varphi + \mu_h + e\varphi = \mu_e + \mu_h = 0$$

under illumination

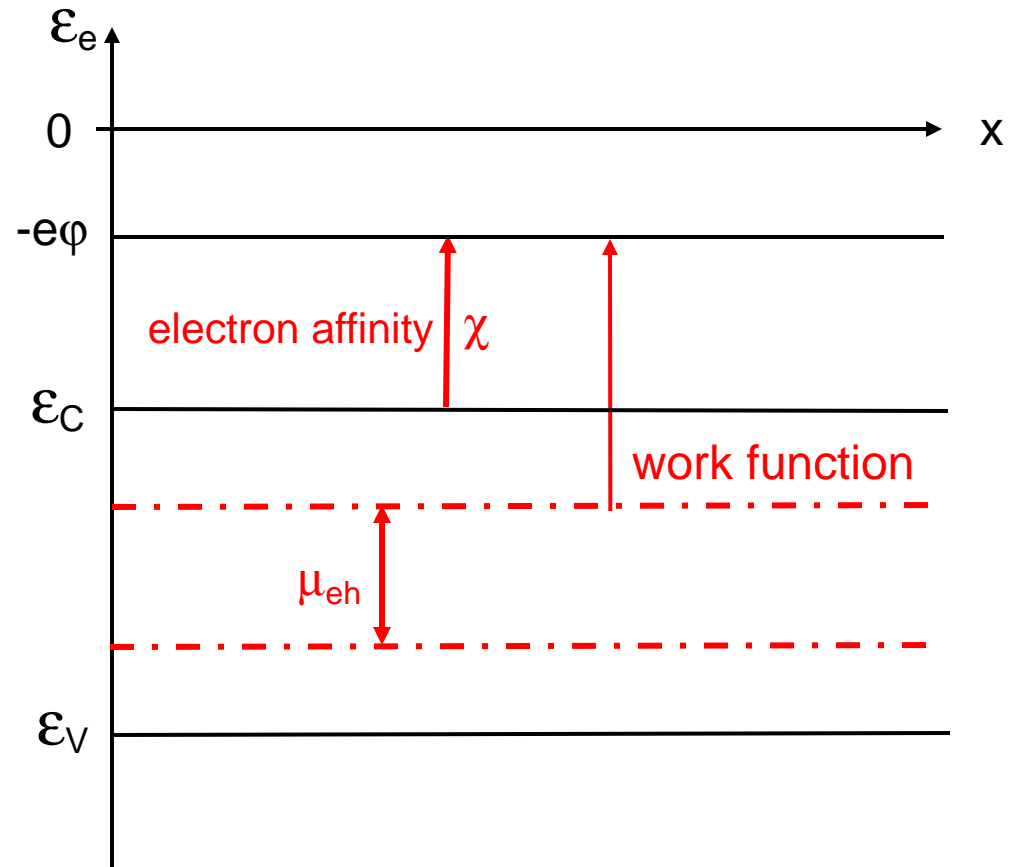
$$\mu_e + \mu_h = \varepsilon_{FC} - \varepsilon_{FV} > 0$$

Free energy of electron-hole pair
at one location is chemical energy

Energy scale for electrons

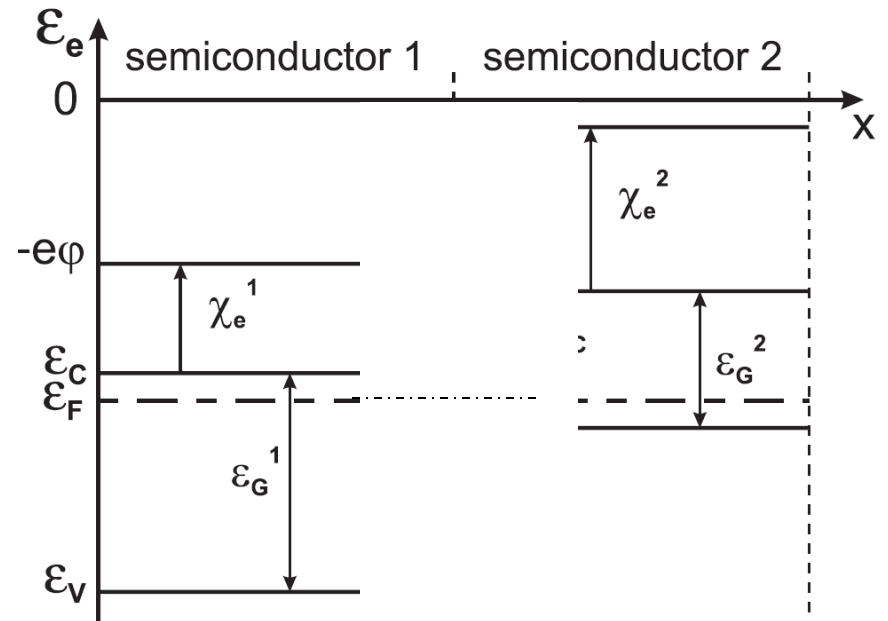
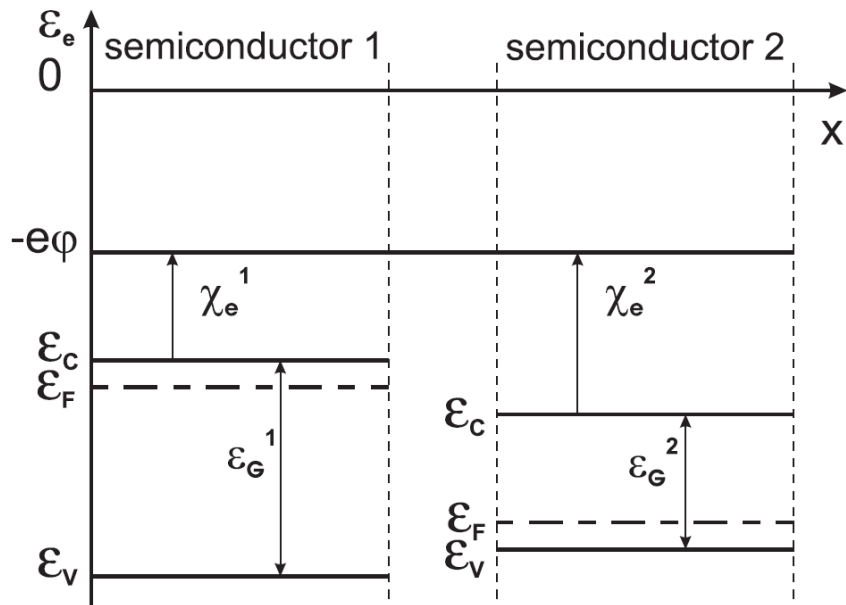
$\varepsilon_e = 0$ for a free electron

1. in vacuum
2. in el. potential $\varphi = 0$
3. with $\varepsilon_{e,\text{kin}} = 0$



energy scale for holes ?

Contact between different materials



Density of states for photons $D_\gamma(\hbar\omega)$

energy distribution of states from uncertainty principle $(\Delta x)^3 (\Delta p)^3 \geq h^3$

volume in momentum space for 1 state $(\Delta p)^3 = \frac{h^3}{V}$

number of momentum states with $|p'| \leq |p|$
i.e. in a sphere with radius $|p|$

$$N(|p|) = \frac{4\pi/3 |p|^3}{h^3/V}$$

$$\varepsilon_\gamma(p) = \hbar\omega = c|p|$$

2 polarisation directions

$$N_\gamma(\hbar\omega) = 2 \times \frac{4\pi V}{3h^3 c^3} (\hbar\omega)^3$$

$$D_\gamma(\hbar\omega) = \frac{1}{V} \frac{dN_\gamma(\hbar\omega)}{d\hbar\omega} = \frac{8\pi}{h^3 c^3} (\hbar\omega)^2$$

in solid angle Ω :
$$D_{\gamma,\Omega}(\hbar\omega) = \frac{2\Omega}{h^3 c^3} (\hbar\omega)^2 = \frac{\Omega}{4\pi^3 \hbar^3 c^3} (\hbar\omega)^2$$

in matter: $c = c_0 / n$

Planck's law

photon current density per solid angle

$$dj_{\gamma,\Omega}(\hbar\omega) = cD_{\gamma,\Omega}(\hbar\omega) f_{\gamma}(\hbar\omega) d\hbar\omega$$

$$D_{\gamma,\Omega}(\hbar\omega) = \frac{2}{h^3 c^3} (\hbar\omega)^2 = \frac{(\hbar\omega)^2}{4\pi^3 \hbar^3 c^3} \quad f_{\gamma}(\hbar\omega) = \frac{1}{\exp\left(\frac{\hbar\omega}{kT}\right) - 1}$$